

Langevin dynamics with multiplicative noise: discretisation issues and path-integral representations

Leticia Cugliandolo⁽¹⁾, Vivien Lecomte⁽²⁾, Frédéric van Wijland⁽³⁾

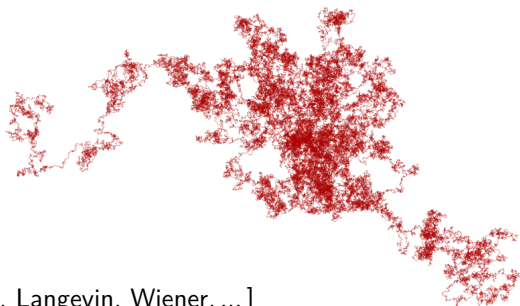
⁽¹⁾LPTHE, Paris

⁽²⁾LIPhy, Grenoble

⁽³⁾MSC, Paris

SISSA / ICTP – Trieste

Fluctuating trajectories – generic settings



[Brown, Einstein, Langevin, Wiener, ...]

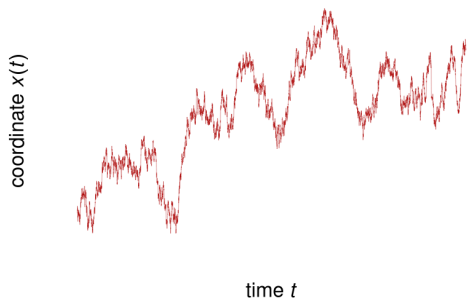
$$\underbrace{m\ddot{x}(t)}_{\text{inertia}} + \underbrace{\gamma\dot{x}(t)}_{\text{damping}} = \underbrace{f(x(t))}_{\text{force}} + \underbrace{\xi(t)}_{\text{noise}}$$

where

$x(t)$ = position, magnetisation, field, ...

$\xi(t)$ = result of from many small contributions (bath, ext. forces, ...)

Fluctuating trajectories – the case studied here



$$\underbrace{\dot{x}(t)}_{\text{damping}} = \underbrace{f(x(t))}_{\text{force}} + \underbrace{g(x(t))\eta(t)}_{\text{multiplicative noise}}$$

The *overdamped* dynamics has:

one dimension

“multiplicative noise” (*i.e.* configuration-dependent noise)

with centered Gaussian $\eta(t)$

$$\langle \eta(t)\eta(t') \rangle = 2D\delta(t - t') \quad (\text{white noise})$$

Statement of the problem

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- Can one do the same for the **trajectory probability**?

$$\mathcal{D}x \mathbb{P}[(x(t))_{0 < t < t_f}] = \mathcal{D}x \exp \left\{ - \int_0^{t_f} dt \mathcal{L}(x, \dot{x}) \right\}$$

Outline

- Introduction
- Discretising Langevin equations
 - Covariance in continuous time (Stratonovich scheme)
 - Covariance in discrete time
- Discretising path-integral trajectory probabilities $\mathbb{P}[(x(t))_{0 < t < t_f}]$
 - The Stratonovich scheme is **not** covariant
 - An “adaptive” covariant discretisation scheme

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What is discretisation?

Discretisation of $[0, t_f]$:

$N = t_f/\Delta t$ time steps of duration Δt

$x_t = x(t)$ at discrete times $t = 0, \Delta t, 2\Delta t, \dots$

Discretisation of Langevin: denoting $\Delta x = x_{t+\Delta t} - x_t$

$$\dot{x}(t) = f(x(t)) + g(x(t))\eta(t)$$



$$\frac{\Delta x}{\Delta t} = f(\bar{x}_t) + g(\bar{x}_t)\eta_t \quad \text{with } x_t \leq \bar{x}_t \leq x_{t+\Delta t}$$

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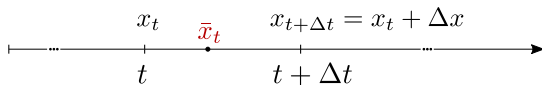


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Noise distribution: η_t 's Gaussian i.i.d. with $\langle \eta_t \eta_{t'} \rangle = 2D\delta_{tt'}/\Delta t$

$$\mathbb{P}(\eta_t) = \left[\frac{\Delta t}{4\pi D} \right]^{\frac{1}{2}} e^{-\frac{\Delta t}{4D}\eta_t^2} \implies \boxed{\eta_t = O(\Delta t^{-1/2}) \implies \Delta x = O(\Delta t^{1/2})}$$

Do we need discretisation?



Discretisation of Langevin: with $\Delta x = x_{t+\Delta t} - x_t$

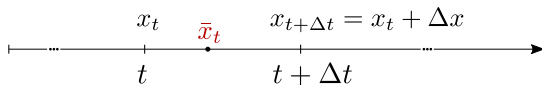
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Scheme: $\bar{x}_t = x_t + \alpha \Delta x$ (with $0 \leq \alpha \leq 1$)

Itô ($\alpha = 0$): $\bar{x}_t = x_t$ (simple for numerics)

Stratonovich ($\alpha = \frac{1}{2}$): $\bar{x}_t = (x_t + x_{t+\Delta t})/2$ (simple for time reversal)

Do we need discretisation? **YES**



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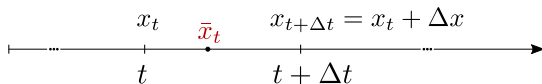
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Fixing α is essential: upon Taylor expansion

$$\frac{\Delta x}{\Delta t} = f(\bar{x}_t) + g(\bar{x}_t)\eta_t = f(x_t) + g(x_t)\eta_t + \underbrace{\alpha \Delta x g'(x_t)}_{=O(\Delta t^0)} \eta_t + O(\Delta t^{\frac{1}{2}})$$

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Conclusion: the continuous-time writing $\dot{x} = f(x) + g(x)\eta$ $[\Delta t \rightarrow 0]$
is *ambiguous* unless one specifies α

A covariant discretisation?

$$u(t) = U(x(t))$$

Naive change of variables: apply the standard **chain rule** $\dot{u} = U'(x) \dot{x}$

$$\dot{u} = F(u) + G(u)\eta$$

with $F(u) = f(x)U'(x)$ and $G(u) = g(x)U'(x)$

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Is there an α -scheme for which it works?

$$\begin{aligned} \Delta u &= U(x_{t+\Delta t}) - U(x_t) = U(x_t + \Delta x) - U(x_t) \\ &= U'(x_t)\Delta x + \underbrace{\frac{1}{2}U''(x_t)\Delta x^2}_{=O(\Delta t)} + O(\Delta x^3) \end{aligned}$$

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In Stratonovich $\bar{x}_t \stackrel{s}{=} x_t + \frac{1}{2}\Delta x$:

$$\frac{\Delta u}{\Delta t} \stackrel{s}{=} U'(\bar{x}_t) \frac{\Delta x}{\Delta t} + O(\Delta t^{1/2})$$

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In Stratonovich:

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An improved scheme: $\bar{x}_t \stackrel{\beta_g}{=} x_t + \frac{1}{2}\Delta x + \beta_g(x_t)\Delta x^2$

$$\beta_g(x) = \frac{1}{24} \frac{g''(x)}{g'(x)} - \frac{1}{12} \frac{g'(x)}{g(x)}$$

improves covariance:

$$\dot{x} \stackrel{\beta_g}{=} f(x) + g(x)\eta$$

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$$\dot{u} \stackrel{\beta_G}{=} F(u) + G(u)\eta + O(\Delta t)$$

For enthusiasts: an **exact** covariant discretisation

Define the $\mathbb{T}_{f,g}$ operator:

$$\mathbb{T}_{f,g}h(x) = \frac{e^{\mathcal{D}(x) \frac{d}{dx}} - \mathbf{1}}{\mathcal{D}(x) \frac{d}{dx}} h(x) = \sum_{n \geq 0} \frac{(\mathcal{D}(x) \frac{d}{dx})^n}{(n+1)!} h(x)$$

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at all orders in Δt

$$\frac{\Delta u}{\Delta t} \stackrel{\mathbb{T}_{F,G}}{=} F(\bar{u}_t) + G(\bar{u}_t)\eta_t$$

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Aims

Path integral: $\langle \mathcal{O}(t_1, \dots, t_k) \rangle = \int \mathcal{D}x O[x(t)] \mathbb{P}[x(t)]$

A very long history:

- Mathematics: PJ Daniell (1919), N Wiener ('20s)
- Quantum mechanics: R Feynman (1948), BS DeWitt (1957)
- Stoch. proc.: Stratonovich (1960), Horsthemke & Bach (1975), Graham (1977), Tirapegui & *al.* (70s-80s)
- Fields: Janssen ('70s), De Dominicis ('70s), Doi & Peliti ('80s)

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Hand-waving procedure:

$$\mathbb{P}[\eta(t)] \propto e^{-\int_0^{t_f} dt \frac{\eta(t)^2}{4D}} \left. \begin{array}{l} \eta = \frac{\dot{x} - f(x)}{g(x)} \\ \dot{x} = f(x) + g(x)\eta \end{array} \right\} \longrightarrow \mathbb{P}[x(t)] \propto \exp \left\{ - \underbrace{\int_0^{t_f} dt \frac{(\dot{x} - f(x))^2}{4Dg(x)^2}}_{\text{action } S[x(t)]} \right\}$$

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 - How to discretise?
 - Can we replace \propto by $=$?

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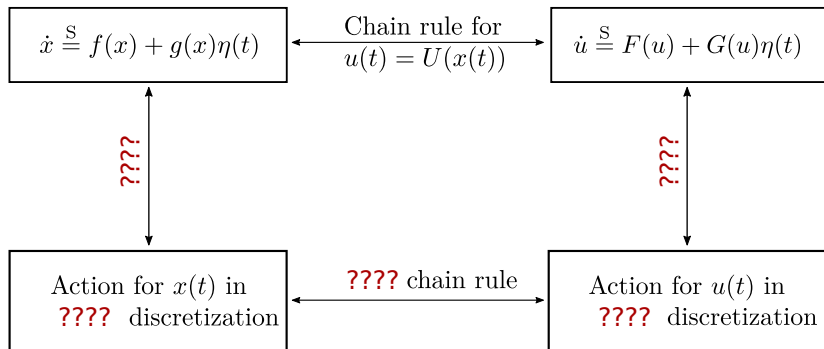
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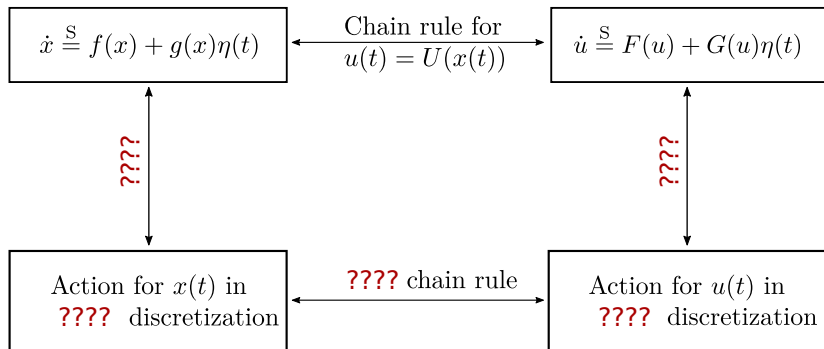
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- Can we make sense of this procedure?
 - How to discretise?
 - Can we replace \propto by $=$?
- Can we ensure **covariance**?
 - The action writes as $S[x] = \int_0^{t_f} dt \mathcal{L}(x, \dot{x})$
 - Can we get $\mathcal{L}(u, \dot{u})$ by applying the **chain rule** in $\mathcal{L}(x, \dot{x})$?

Covariance

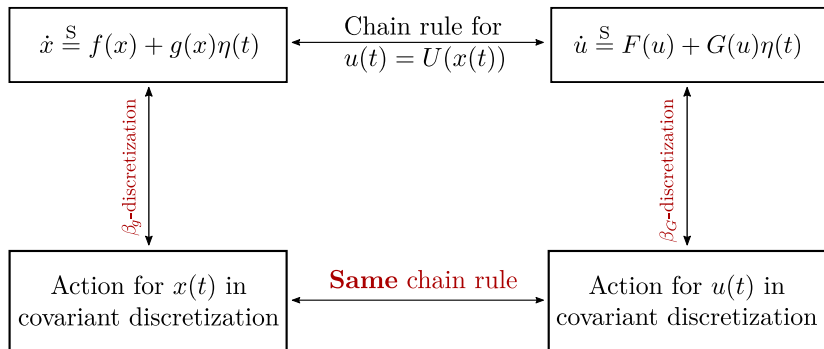


Covariance



Remark: the naive action $S[x(t)] = \int_0^{t_f} dt \frac{(\dot{x} - f(x))^2}{4Dg(x)^2}$ is **not** covariant

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From noise to path trajectory probability

Noise trajectory probability:

$$\mathbb{P}[\{\eta.\}] = \prod_{0 \leq t < t_f} \left[\frac{\Delta t}{4\pi D} \right]^{\frac{1}{2}} e^{-\frac{\Delta t}{4D} \eta_t^2}$$

Recursion at each time step:

$$\frac{x_{t+\Delta t} - x_t}{\Delta t} = f(\bar{x}_t) + g(\bar{x}_t)\eta_t \quad \Rightarrow \quad x_{t+\Delta t} = x_{t+\Delta t}(x_t, \eta_t)$$

Trajectory probability as a Feynman path integral

$$\mathbb{P}[\{x.\}] = \prod_{0 \leq t < t_f} \mathbb{P}(x_{t+\Delta t}, t + \Delta t | x_t, t)$$

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$$\mathbb{P}[\{x.\}] = \prod_{0 \leq t < t_f} \mathbb{P}(x_{t+\Delta t}, t + \Delta t | x_t, t) \equiv \mathcal{N}[\{x.\}] e^{-S[\{x.\}]}$$

$\mathcal{N}[\{x.\}] =$ Normalisation prefactor

$S[\{x.\}] =$ Discrete-time action

A covariant normalisation prefactor

End-point discretised $\mathcal{N}[x(t)]$:

$$\mathcal{N}[\{x.\}] = \prod_{0 \leq k < t_f} \frac{1}{\sqrt{4\pi D \Delta t}} \frac{1}{|g(x_{t+\Delta t})|}$$

Origin of the covariance: compensation of $U'(x_{t+\Delta t})$

- upon changing variables:

$$\mathbb{P}_X(x_{t+\Delta t}|x_t) = |U'(x_{t+\Delta t})| \mathbb{P}_U(u_{t+\Delta t}|u_t)$$

- upon changing between g and G :

$$g(x_{t+\Delta t}) = G(u_{t+\Delta t})/U'(x_{t+\Delta t})$$

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Conclusion:

this choice for $\mathcal{N}[\{x.\}]$ allows one to focus on the sole action

A Stratonovich-discretised action?

Using standard procedures:

$$\mathbb{P}(x_{t+\Delta t}|x_t) \stackrel{s}{=} \frac{1}{\sqrt{4\pi D\Delta t} |g(x_{t+\Delta t})|} e^{-\delta S(\bar{x}_t)}$$

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$$\begin{aligned} \delta S(\bar{x}) \stackrel{s}{=} & \frac{1}{2} \frac{\Delta t}{2D} \left[\frac{\frac{\Delta x}{\Delta t} - f(\bar{x})}{g(\bar{x})} \right]^2 + \frac{\Delta t}{2} \left[f'(\bar{x}) - \frac{f(\bar{x})g'(\bar{x})}{g(\bar{x})} \right] \\ & + \frac{D}{4} [2g'(\bar{x})^2 - g(\bar{x})g''(\bar{x})] \Delta t \end{aligned}$$

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$$\frac{\Delta t}{2Dg(\bar{x}_0)^2} \frac{\Delta x}{\Delta t} \times O(\Delta t^{1/2})$$

which is of order Δt and thus **cannot be neglected**.

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Remark:

- The writing is *different* from that of the Stratonovich discretisation
- The action is thus more sensitive to discretisation details than the Langevin equation [Gulyaev & Edwards 1964, Tirapegui & al. 70s, ...]

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- In the end, a mere **expansion in powers of Δt**
 → See *Building a path-integral calculus* [arXiv:1806.09486]
 Leticia Cugliandolo, VL, Frédéric van Wijland

Remark 1/2

Mathematician's view of Stochastic Calculus:

Give a well-defined meaning to

$$\int_0^{t_f} dt \left[h_1(x(t)) + h_2(x(t)) \dot{x} \right]$$

i.e. go beyond Riemann sum (use a finer discretisation)

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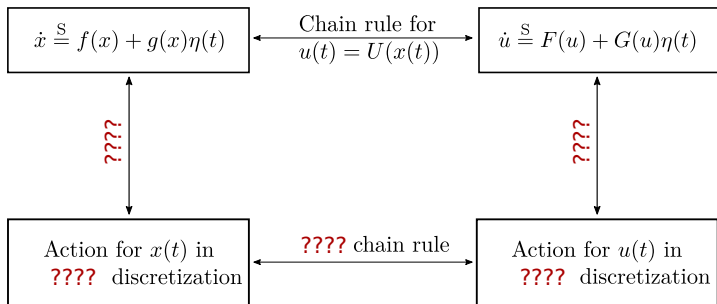
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What the covariant discretisation brings:

Allow for the use of ordinary calculus in

$$\exp \left\{ - \int_0^{t_f} dt \left[h_1(x(t)) + h_2(x(t)) \dot{x} + h_3(x(t)) \dot{x}^2 \right] \right\}$$

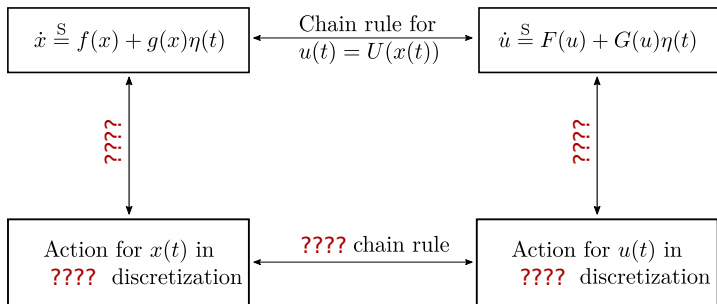
Remark 2/2



For an arbitrary $????$ discretisation:

- The action is still correct
- The use of the ordinary chain rule in the action yields *wrong* results

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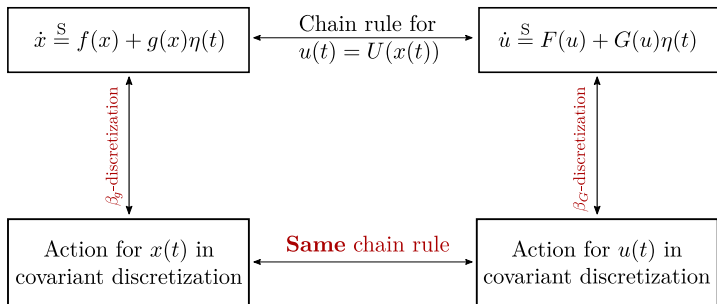
For an arbitrary **????** discretisation:

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- Changing variables in the time-discrete action is still possible

But: yields intricate rules in the continuous-time limit

See Leticia Cugliandolo & VL, JPhysA **50** 345001 (2017)

A graphical summary



An unequivocal meaning to $\mathbb{P}[x(t)] = \mathcal{N}[x(t)]e^{-S[x(t)]}$ and to

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“Adaptive” covariant discretisation:

$$\bar{x}_t \stackrel{\beta_g}{=} x_t + \frac{1}{2}\Delta x + \beta_g(x_t)\Delta x^2$$

Perspectives

Extensions:

- Martin–Siggia–Rose–Janssen–de Dominicis action
- Higher dimensions; **non-trivial!** [w.i.p. Thibaut Arnoux de Pirey]

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- Quantum-mechanical action
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- Supersymmetry for multiplicative noises
- Covariant WKB (small-noise expansion)

Supplementary material

Operator approach

From Fokker–Planck operator to path integrals

$$\begin{aligned}\partial_t P(x, t) &= \mathbb{W} P(x, t) \\ P(x_f, t_f | x_i, t_i) &= \langle x_f | e^{(t_f - t_i) \mathbb{W}} | x_i \rangle \\ &= \int_{x(t_i) = x_i}^{x(t_f) = x_f} \mathcal{D}x O[x(t)] \mathbb{P}[x(t)]\end{aligned}$$